

# ~~4~~-dimensional gauge theory and 3 representation theory

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In Nov. 1996, Witten gave a series of lectures at  
Newton Institute, Cambridge.

He explained 3 topics:

1. 3-manifolds invariants via hyper-Kaehler manifolds  
(Rozansky-Witten)
2. Coulomb branch for 3d  $N=4$  SUSY gauge theory  
(Seiberg-Witten for  $SU(2)$  and others)
3. 3d mirror symmetry (Intriligator-Seiberg) and  
its description by branes (Hanany-Witten)

They were instructive lectures, and I learned quite lots.

But I could not produce my own work out of them at that time.

Today : My response to lectures, 18 years later !

## 3d N=4 SUSY gauge theory (topologically twisted version)

### data

- $G$  : a compact Lie group
- $V$  : a quaternionic representation of  $G$   $(G \times Sp(1) \cong Spin(3))$
- $P$  : a principal  $G$ -bundle over  $M^3$

### fields

- $A$  :  $G$ -connection + etc (vector multiplet)
- $\Phi$  : spinor for  $(P \times P_{Spin})_{G \times Sp(1)} \times V$  + etc (hypermultiplet)

$\rightsquigarrow$  a hyper-Kaehler manifold  $\mathcal{M}_C$  : Coulomb branch

$\mathcal{M}_C$  is important to understand the gauge theoretic 3d invariants, defined by counting solutions of PDE of fields.

e.g. Casson – Walker invariant  $(G = SU(2), V = 0)$  = Rozansky – Witten invariant of  $\mathcal{M}_C$   
 3d monopole invariant  $(G = U(1), V = \mathbb{H})$

Seiberg-Witten (1996)

$$G = SU(2) \ , \ T = 0 \quad \Rightarrow \quad \mathcal{M}_C = \text{Atiyah-Hitchin manifold}$$

$$G = U(1) \ , \ T = \mathbb{H} \quad \Rightarrow \quad \mathcal{M}_C = \text{Taub-NUT space}$$

$\mathcal{M}_C$  is based on quantum field theory, which lacks mathematical foundation. For mathematicians, claims sound as

(something, not defined) = AH, or TN space.

Thus mathematicians cannot understand these statements.

Later  $\mathcal{M}_C$  were computed in many examples in physics literature, interesting hyper-Kaehler manifolds, such as instanton moduli spaces of exceptional groups, appeared.

Examples.

## 1. Abelian gauge theory

$$G = U(1)^n \hookrightarrow \mathbb{H}^d$$

$$\implies 0 \rightarrow \mathbb{Z}^n = \pi_1(G) \longrightarrow \mathbb{Z}^d = \pi_1(U(1)^d) \longrightarrow \mathbb{Z}^{d-n} \rightarrow 0$$

Apply  $\text{Hom}(\cdot, U(1))$  :

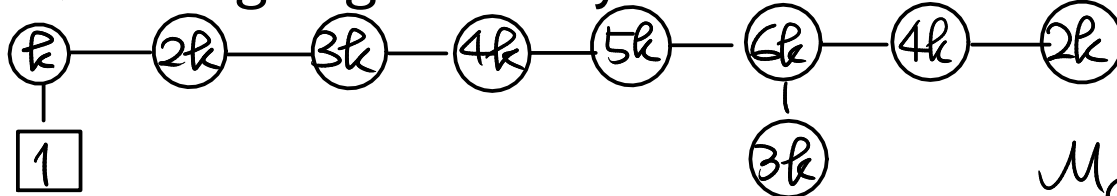
$$1 \rightarrow U(1)^{d-n} \longrightarrow U(1)^d \longrightarrow G^V (\text{dual torus}) \rightarrow 1$$

Then  $\mathcal{M}_C = (\mathbb{H}^d)^* // U(1)^{d-n}$  (hyperkähler quotient)

$$\text{cf. } \mathcal{M}_H = \mathbb{H}^d // G$$

## 2. Quiver gauge theory

( $k \in \mathbb{Z}_{>0}$ )



$\mathcal{M}_C =$  moduli space of  $k$   $E_8$ -instantons on  $\mathbb{R}^4$

It was similar to Seiberg-Witten curves for 4d N=2 SUSY gauge theory. Analogue of  $\mathcal{M}_C$  was determined as a family of elliptic curves (Seiberg-Witten curves).

A mathematical definition of (analogue of)  $\mathcal{M}_C$  was given by Nekrasov much later in 2002 by  $\Omega$ -background.

And the theory has been developed further after that.

I learned from this lesson:

Writing down a mathematically rigorous definition is a non-trivial, often difficult, but challenging problem.

Today : I propose a mathematically rigorous definition of  $\mathcal{M}_C$ ,  
 at least as a complex manifold (for good or ugly theories).  
 (work in progress with Braverman+Finkelberg)

$$\Sigma = \mathbb{P}^1$$

$\mathcal{R}$  = moduli stack of pairs

- $\mathcal{E}$  : a holomorphic  $G_{\mathbb{C}}$ -bundles over  $\Sigma$
- $\Phi$  : a holomorphic section of  $(\mathcal{E} \times_{G_{\mathbb{C}}} V)^{\otimes \frac{1}{2}} \otimes \mathcal{O}_{\Sigma}(-1)$  such that

$$\mu_{\mathbb{C}}(\Phi) = 0$$

$\mathcal{R}$  is a critical point of a complex Chern-Simons type functional,  
 hence has a sheaf  $\mathcal{V}_{CS}$  of the vanishing cycle.

$$\mathbb{A}[\mathcal{M}_C] = (\text{dual of}) \ H_c^*(\mathcal{R}, \mathcal{V}_{CS})$$

$\uparrow$   
 $U(1) \subset SU(2)$

$\uparrow$   
 cohomology degree

(Multiplication will be explained later.)

Reasons why this definition should give the Coulomb branch  $\mathcal{M}_C$

1. It reproduces the monopole formula of Cremonesi-Hanany-Zaffaroni, found recently (at least if  $\nabla = \nabla' \oplus_{\mathcal{H}_G} (\nabla')^*$  ).
2. When  $V = 0$  (which violates good or ugly assumption), Bezrukavnikov-Finkelberg(-Mirkovic) identified  $\mathcal{M}_C$  with a certain symplectic quotient, which is moduli space of charge  $k$   $SU(2)$ -monopoles (for  $SU(k)$ ) (Bielawski), as predicted by Seiberg-Witten ( $k=2$ ) and Chalmers-Hanany.
3. True for abelian
4. A natural quantization



1.  $G_{\mathbb{C}}$  - bundles over  $\mathbb{P}^1$  are reduced to  $T_{\mathbb{C}}$ -bundles.  
 (Vector bundles over  $\mathbb{P}^1$  are direct sum of line bundles.)

This induces a stratification of  $\mathcal{R}$ , which is \*perfect\*.

Therefore the dimension of the cohomology is the sum of dimension of cohomology of strata.

$$P_t(H_c^*(\mathcal{R}, \mathcal{U}_{CS})) = \sum_{\substack{\lambda \in \text{coweight of } G \\ \text{Weyl group}}} \frac{1}{P_t(H_{\text{Stab}(\lambda)}^*(pt))} \times t^{2\Delta(\lambda)}$$

$$\Delta(\lambda) = - \sum_{\alpha \in \Delta^+} |\langle \alpha, \lambda \rangle| + \frac{1}{2} \sum_{\substack{b: \text{H-base} \\ \text{of } T}} |\langle \omega + b, \lambda \rangle|$$

Product

Taking the Poincare dual, we switch from bundles over  $\mathbb{P}^1$  to affine Grassmannian  $\text{Gr}_{\mathbb{G}_{\mathbb{C}}}$   
= bundles over  $\mathbb{P}^1$  + trivialization away  $\infty$

$\text{Gr}_{\mathbb{G}_{\mathbb{C}}} \sim \Omega G$  : based loops group

$\implies$  Pontryagin product  $H_*(\Omega G) \otimes H_*(\Omega G) \longrightarrow H_*(\Omega G)$